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MEASUREMENT OF THE REAL PART OF THE
FORWARD ELASTIC NUCLEAR AMPLITUDE
FOR K^-p AT 16 GeV/c

by

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ABSTRACT

The differential cross section for the K^-p elastic scattering at 16 GeV/c in the range $0 < |t| \leq 0.4 \text{ GeV}^2$ is obtained. In the t -range up to 0.1 GeV^2 the ratio ρ of the real to the imaginary part of the nuclear amplitude and the slope parameter are measured. For the other t -range a conventional exponential fit is used in order to determine the slope parameter. A comparison between ρ and theoretical predictions from dispersion relations is made.

1. INTRODUCTION

We have investigated the elastic scattering of K^- mesons on protons for small momentum transfers ($0.01 < |t| < 0.4$) GeV^2 . We have taken into account the Coulomb amplitude. It is comparable in strength to the nuclear amplitude in the range ($0.01 < |t| < 0.1$) GeV^2 , giving rise to interference effects. In this way we could determine both the magnitude and sign of the real part of the scattering amplitude, at small values of t .

Through the determination of the ratio ρ of the real to the imaginary part of the forward elastic amplitudes for different incident energies^[1-10], the validity of forward dispersion relations can be checked^[11-14].

This paper presents our determination of ρ for K^-p at 16 GeV/c (where, up to now, there is no information) from a sample of 41.332 elastic events^(*) observed in the CERN 2m hydrogen bubble chamber.

In Sect. 2 we describe the method that we used to treat the data sample at very small values of t . In Sect. 3 we give the differential cross section and we discuss the parametrization used to determine ρ . Sect. 4 contains the comparison between our experimental result with the dispersion relation curve.

(*) We thank Prof. G. Otter and the members of the Aachen-Berlin-CERN-London-Vienna Collaboration who made these data available to us.

2. THE TREATMENT OF THE DATA

In order to correct for scanning and measurement losses of events at small values of t (up to 0.2 GeV^2), we have made a scatter diagram for each t interval plotting the momentum distribution of the scattered K^- mesons in the (p_y, p_z) plane perpendicular to the beam direction. In this way, all the events with a momentum p_y smaller than a limiting value p_c (cut-off) has been geometrically corrected on basis of cylindrical symmetry around the beam direction.

Fig. 1 shows an example of such a scatter diagram for the $|t|$ interval 0.040 to 0.045 GeV^2 .

3. PARAMETRIZATION OF THE DIFFERENTIAL CROSS SECTION

The differential cross-section $d\sigma/dt$ (Fig. 2) has been obtained by normalization to the observed value of the integrated elastic cross section $\sigma_{e1} = 2.77 \pm 0.17 \text{ mb}$ [15].

Following ref. [3], we have taken the Coulomb amplitude as

$$F_C(t) = 2\alpha\hbar\sqrt{\pi} \beta_L^{-1} e^{2i\delta} e^{-\frac{1}{2}b|t|} |t|^{-1} \quad (1)$$

where α is the fine structure constant, β_L is the velocity of the K^- meson in the laboratory system, b is the slope parameter (the same for the nuclear amplitude) and δ is the model-dependent phase of the Coulomb amplitude [16] given by

$$\delta = -(\ln|t| + \ln b + 0.577)\alpha(2\beta_L)^{-1} \quad (2)$$

The nuclear amplitude is taken as

$$F_N(t) = \sigma_t (4\sqrt{\pi} \hbar c)^{-1} (\rho+i) e^{-\frac{1}{2}b|t|} \quad (3)$$

where σ_t is the total K^-p cross section at 16 GeV/c and

$$\rho = \text{Re } F_N(0) / \text{Im } F_N(0) \quad (4)$$

σ_t is related to the imaginary part of the forward nuclear amplitude through the optical theorem.

The differential cross section is given by

$$\frac{d\sigma}{dt} = |F_C + F_N|^2 = \left(\frac{d\sigma}{dt}\right)_C + \left(\frac{d\sigma}{dt}\right)_N + \left(\frac{d\sigma}{dt}\right)_I \quad (5)$$

where C, N and I represents Coulomb, Nuclear and Interference terms. Taking F_C and F_N as defined above we have

$$\left(\frac{d\sigma}{dt}\right)_C = 2.60945 \times 10^{-4} t^{-2} e^{-bt} \quad (6)$$

$$\left(\frac{d\sigma}{dt}\right)_N = 23.8820 (1+\rho^2) e^{-bt} \quad (7)$$

and

$$\left(\frac{d\sigma}{dt}\right)_I = 0.157884 t^{-1} (\rho \cos 2\delta - \sin 2\delta) e^{-bt} \quad (8)$$

We have determined values of the parameters ρ and b (assuming them to be independent) by making a least squares fit in the range $(0.01 < |t| < 0.1)\text{GeV}^2$. We have obtained the values

$$\rho = 0.123 \pm 0.096$$

and

$$b = 9.50 \pm 0.65 \text{ GeV}^{-2} .$$

The value of the differential cross section for the first t bin was not taken in the fit due the fact that the scanning loss corrections are not reliable.

For $0.1 < t \leq 0.4$, the contribution of the Coulomb scattering is small and can be neglected. We have fitted the data in this region by a simple exponential e^{-Bt} . We have found the value

$$B = 7.04 \pm 0.23 \text{ GeV}^{-2} .$$

In Fig. 3 a and b we show the curve of our fit for these two regions of t , respectively.

4. CONCLUSIONS

As we can see in Fig. 4, our result for ρ agrees well with the prediction given by dispersion relations [17,18].

We have found a change of slope in the $d\sigma/dt$ distribution around $|t| \sim 0.1 \text{ GeV}^2$. This was also observed by Campbell et al. at 10.1 GeV/c.

ACKNOWLEDGEMENTS

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FIGURE CAPTIONS

FIG. 1. - Scatter diagram of the momentum distribution of the scattered K^- mesons in the (p_y, p_z) plane for the t -bin 0.040 to 0.045 GeV^2 .

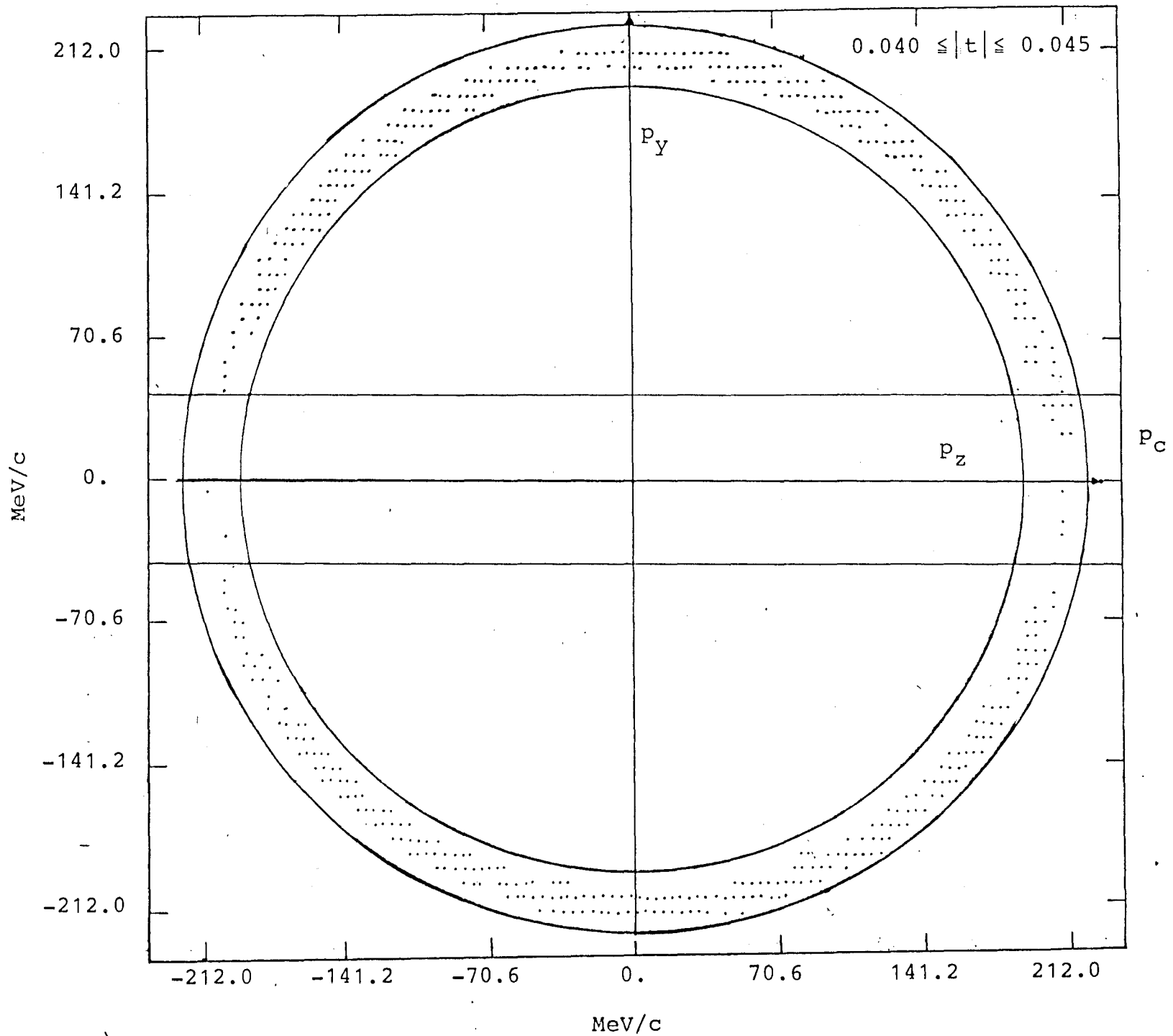
FIG. 2. - K^-p elastic differential cross section at 16 GeV/c after scanning and measurement corrections.

FIG. 3.a - K^-p elastic differential cross section at 16 GeV/c for $0 < |t| \leq 0.1 \text{ GeV}^2$. The line is our fit with $\rho = 0.123 \pm 0.096$ and $b = 9.50 \pm 0.65 \text{ GeV}^{-2}$.

3.b - K^-p elastic differential cross section at 16 GeV/c for $0.1 < |t| \leq 0.4 \text{ GeV}^2$. The line is our fit with $B = 7.04 \pm 0.23 \text{ GeV}^{-2}$.

FIG. 4. - K^-p ratio of forward real part to imaginary part. Experimental points are taken from the following references: \diamond our result; \times Ref. [10]; \triangle Ref. [4]; \circ Ref. [8]; \blacktriangle Ref. [6]; \bullet Ref. [9]; \square Ref. [1]; \blacksquare Ref. [3]; \blacklozenge Ref. [5]. The full line is the theoretical prediction from dispersion relations Ref. [17] and the dashed line Ref. [18].

Fig. 1



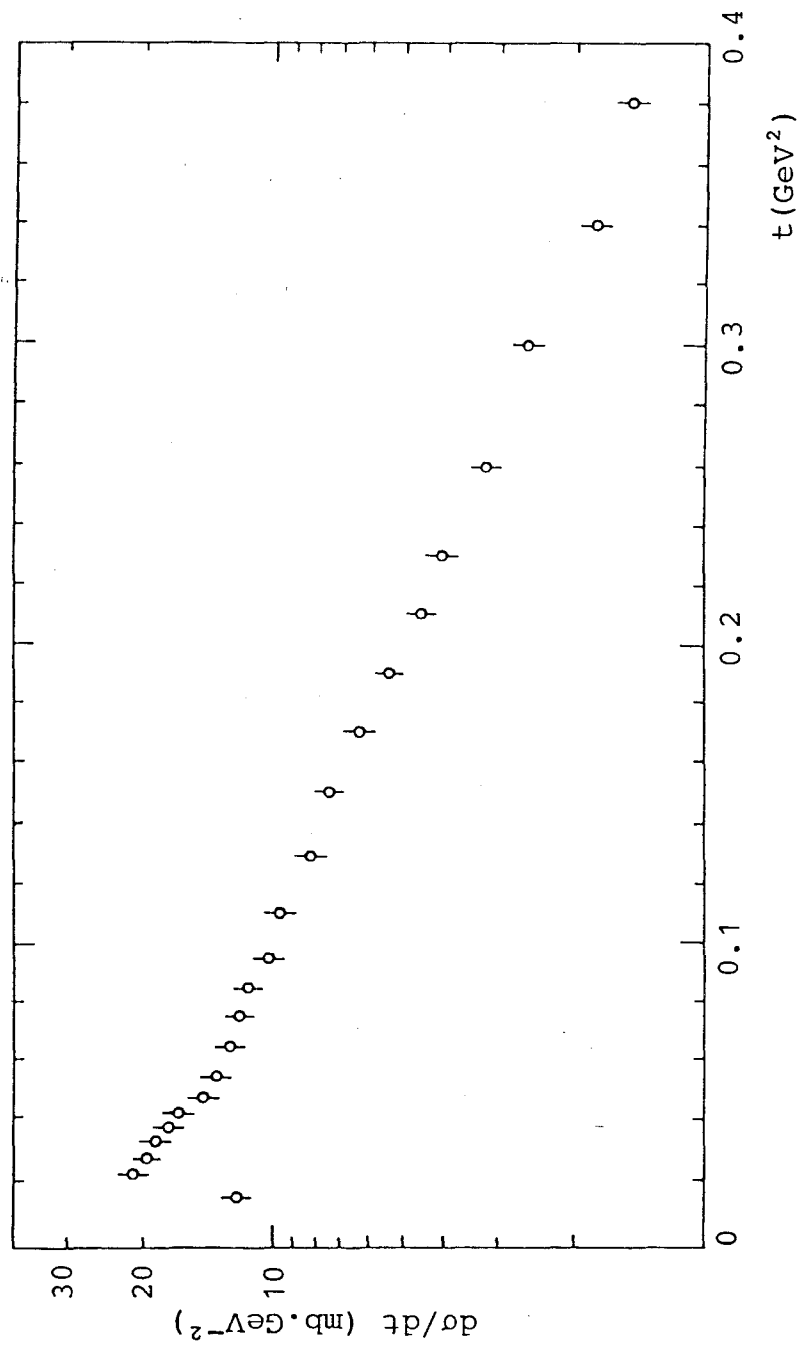


Fig. 2

Fig. 3.a

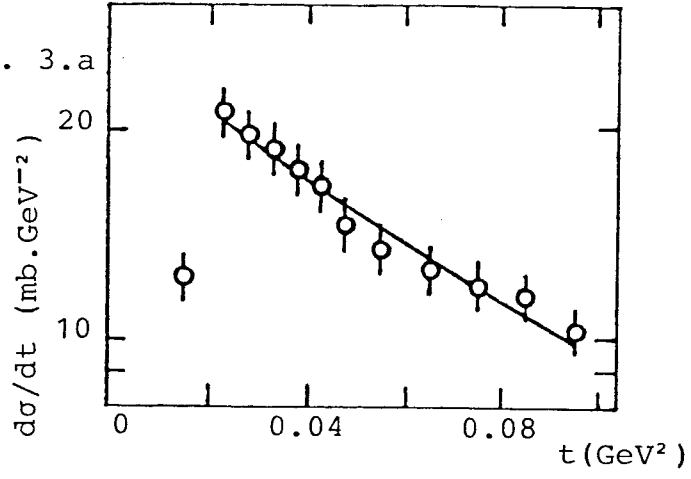


Fig. 3.b

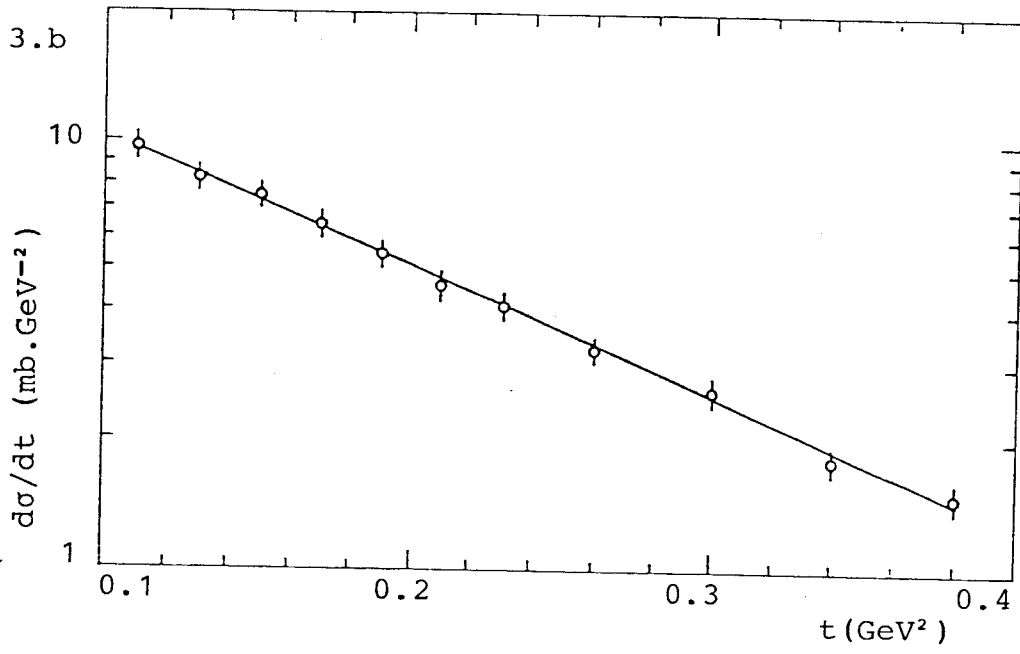


Fig.4

