

CBPF-NF-017/81

ELECTRICAL RESISTIVITY OF CeFe_2

by

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ABSTRACT

Electrical and magnetic properties of CeFe_2 were investigated from the temperature dependence of electrical resistivity in the range of 1,5 to 300°K. The critical temperature, determined from the maximum of $d\rho/dT$, gives $T_c = 220^\circ\text{K}$ and the temperature independent magnetic resistivity is $87\mu\Omega\text{ cm}$. This value is compared with the corresponding in YFe_2 . At low temperature the resistivity shows a fairly large variation proportional to AT^2 up to about 32°K with $A = 1,2 \times 10^{-2} \mu\Omega\text{ cm}/^\circ\text{K}^2$.

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Abstract. Electrical and magnetic properties of CeFe_2 were investigated from the temperature dependence of electrical resistivity in the range of 1,5 to 300°K. The critical temperature, determined from the maximum of $d\rho/dT$, gives $T_c = 220^\circ\text{K}$ and the temperature independent magnetic resistivity is $87\mu\Omega\text{cm}$. This value is compared with the corresponding in YFe_2 . At low temperature the resistivity shows a fairly large variation proportional to AT^2 up to about 32°K with $A = 1,2 \times 10^{-2} \mu\Omega\text{cm}/^\circ\text{K}^2$.

Intermetallic alloys AFe_2 has been investigated very little from transport phenomena measurements. The electrical resistivity for A transition metal as Sc, Y, Ti, Zr, Hf, Nb, Ta (Ikeda et al 1975) and Ti (Ikeda et al 1972) was reported and also for A rare-earth as Tb (Savage 1974). As the short-range ($T > T_c$) and long-range ($T < T_c$) order affects strongly the temperature dependence of the electrical resistivity it provides a good means to study magnetic interactions. However, the measurements of rare-earth compounds are specially hard for the brittleness of the samples.

The compound CeFe_2 , cubic with MgCu_2 type Laves phase is ferromagnetic and has the iron moment and ordering temperatu

re lower than others rare-earth-Fe₂ compounds. Also, in CeFe₂ Ce is quadrivalent and the itinerant electron concentration must be greater. In this paper electrical and magnetic properties of CeFe₂ are investigated from results of electrical resistivity measurements.

The sample preparation follows the usual arc furnace cast (1,5 mm diameter and 1,2 cm long), of elements in stoichiometric composition followed by annealing at 700°C in an argon atmosphere during one week. X-ray using CuK_α radiation shows a single phase with a lattice parameter of 7,302 Å, derived by least squares analysis using Nelson-Riley's extrapolation. This reduced value is due to the +4 valence of Ce.

Electrical resistivity was measured by DC four-points probe method and over a temperature range of 1.5-300°K measured with Ge and Pt sensors. Pressure contacts are used. No thermal hysteresis was observed. Uncertainties in values of ρ are estimated to be about 1% max. due mainly to difficulties in measuring sample dimensions.

The residual resistivity ρ_0 of 4,19 $\mu\Omega$ cm is consistent with that observed in YFe₂ (Ikeda et al 1975) and others AFe₂.

In Fig. 1 the resistivity ρ versus temperature T curve and the plot of $d\rho/dT$ versus T, computed point by point in a three point span to a fit of third-order-polynomial show a ferromagnetic behavior. A sharp break in the slope of ρ vs T at 230°K is observed. The Curie temperature determined from magnetization data range from 230 to 237°K (Buschow et al 1970, Farrell et al 1964, da Cunha et al 1980). But the derivative curve shows a maximum at T = 220°K, the "resistivity

Curie temperature", this difference of T_c been observed for others Laves phases (Kawatra et al 1970).

From low temperature $d\rho/dT$ increases linearly with T until 270K, followed by a broad maximum. Then $d\rho/dT$ decreases nearly until the vicinity of the transition region. This decrease reflects the negative curvature of ρ vs T curve in this range of temperature, and shows clearly that another scattering mechanism takes place. There is very little theoretical investigations for this behavior of ρ versus T .

The linear part in ρ above T_c gives $d\rho/dT = 9 \times 10^{-2} \mu \Omega \text{ cm}/^\circ\text{K}$ value some higher than that of YFe_2 (3×10^{-2}) (Ikeda et al 1975). It should be said that such a value was obtained from measurements made only 70 degrees above T_c , as compared to about 300 degrees in the case of YFe_2 . As the short-range order participates of the scattering mechanism, we need to know δ , the extend of the short-range order. If we take the tail in the $d\rho/dT$ above T_c as a measure of δ this gives ~ 25 degrees.

Assuming the Mathiessen's rule the magnetic resistivity ρ_m due to spin disorder was estimated by extrapolating the linear high temperature part of the resistivity curve to $T=0$ and subtracting the residual resistivity. This yield a value of $87 \mu \Omega \text{ cm}$ which is of the same order of magnitude as that of others AFe_2 . A check of this value can be estimated in the following way.

In the isostructural compound YFe_2 we have a similar electronic structure (few d electrons in the A site) and nearly the same ρ vs T behavior (Ikeda et al 1975). Although CeFe_2 is itinerant, at high temperatures ($T \gg T_c$) we can apply (Ueda

et al 1975), the spin disorder resistivity for ferromagnetic metals given by the Kasuya (1956) expression. Comparing ρm_1 for $CeFe_2$ with ρm_2 for YFe_2 we have:

$$\frac{\rho m_1}{\rho m_2} = \frac{A_1}{A_2} \left(\frac{J_1}{J_2} \right)^2 \frac{S_1(S_1+1)}{S_2(S_2+1)}$$

where

$$A = \frac{3\pi m^*}{e^2 \hbar^2} \frac{1}{NE_F}$$

J is the s-d exchange integral and $S(S+1) = \langle s^2 \rangle$ is the iron's spin.

We can suppose $J_1 \approx J_2$ and $m_1^* \approx m_2^*$.

Then

$$\frac{\rho m_1}{\rho m_2} = \frac{N_2}{N_1} \frac{E_{F_2}}{E_{F_1}} \frac{S_1(S_1+1)}{S_2(S_2+1)}$$

Taking $S_1(S_1+1) = 0,96$, $S_2(S_2+1) = 1,226$ and $N = 24/a^3 \text{ \AA}^{-3}$, $a_1 = 7,302 \text{ \AA}$ for $CeFe_2$ and $a_2 = 7,359 \text{ \AA}$ for YFe_2 (da Cunha et al 1980) we have:

$$\frac{\rho m_1}{\rho m_2} = 0,8 \frac{E_{F_2}}{E_{F_1}}$$

Measurements made by Ikeda and al (1975) in YFe_2 gives a value of the $\rho m_2 = 140 \mu \Omega \text{ cm}$. Then

$$E_{F_1} = 1,3 E_{F_2}$$

As is known that Ce is quadrivalent in CeFe_2 by transferring a f electron to the d-band of the compound is reasonable that Fermi energy in CeFe_2 be larger than in YFe_2 . The iron magnetic moment in YFe_2 is larger than in CeFe_2 . Assuming a common density of states curve for both YFe_2 and CeFe_2 , the f electron transfer to the d-band amounts to displacing the Fermi level to higher energies. If one has a decreasing density of states in this region, the Fe magnetic moments can be also consistent with the f electron transfer.

In Fig. 2 the logarithmic plot of $\rho - \rho_0$ versus T shows a T^2 dependence up to 320K. The best computer fitting of the low temperature curve ρ versus T uses the expression

$$\rho - \rho_0 = AT^2 + BT^5$$

giving also a predominating term in T^2 up to 320K with a coefficient $A = 1,2 \times 10^{-2} \mu\Omega \text{ cm}/\text{K}^2$ and $B = 2 \times 10^{-7} \mu\Omega \text{ cm}/\text{K}^5$. The high value of the coefficient A only compares with those obtained in strong paramagnets like YCo_2 or in weak ferromagnets like ZrZn_2 and Sc_3In where A is approximately equal to $5 \times 10^{-2} \mu\Omega \text{ cm}/\text{K}^2$ (Moriya 1979). Then, remains the problem of explaining this order of magnitude of coefficient A in CeFe_2 .

Acknowledge - The authors wish to acknowledge A.A. Gomes for many fruitful discussions.

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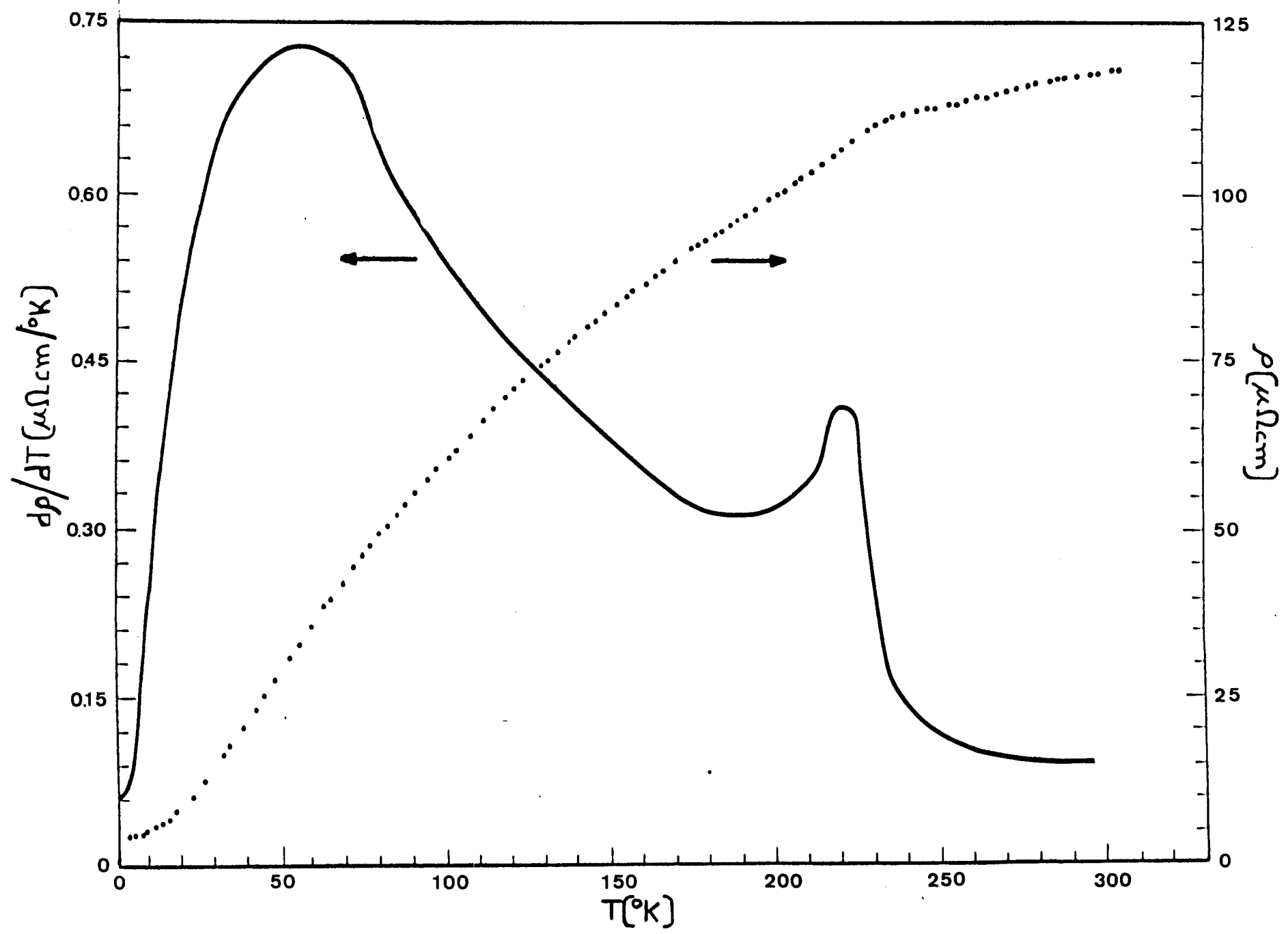
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CAPTIONS

Figure 1 - The temperature variation of the electrical resistivity $\rho(T)$ and of its coefficient $d\rho/dT$.

Figure 2 - Logarithmic plot of $\rho - \rho_0$ versus T at low temperature.

Fig. 1



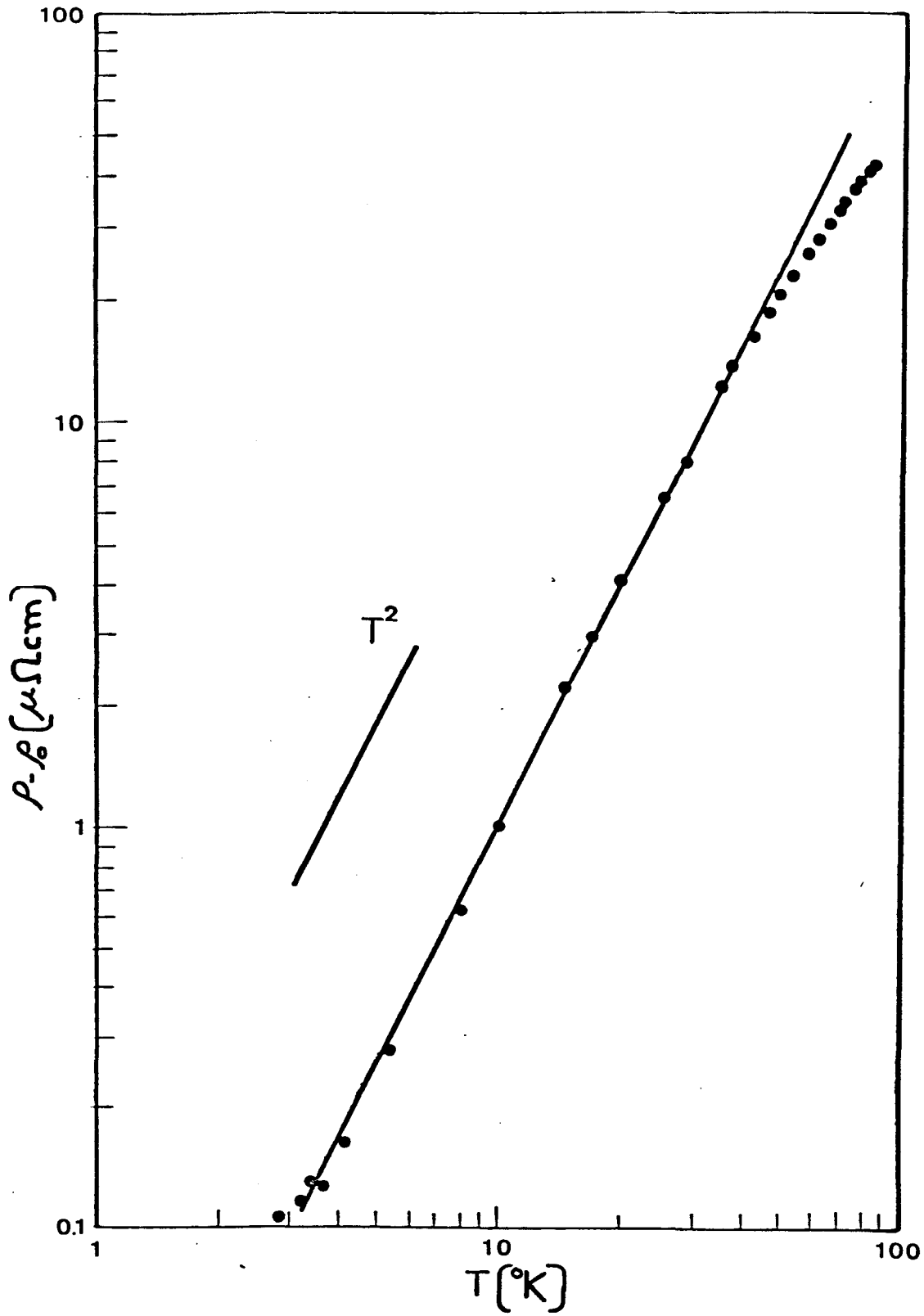


Fig. 2