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ANHYDROUS COBALTOUS CHLORIDE

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ABSTRACT

The Mössbauer spectra of  $^{57}\text{Co}-\text{CoCl}_2$  reveal  $^{57}\text{Fe}$  in the +2 charge state. Isomer shift and quadrupole splitting are discussed. An interpretation is suggested for the observation of only the ferrous ionic state.

\* \* \*

We report in this paper precision measurements of the Mössbauer spectra of  $^{57}\text{Co}$  doped anhydrous  $\text{CoCl}_2$  sources, between 80 and  $500^\circ\text{K}$ <sup>1</sup>. The main purpose of the work was to get further insight in the formation process of multiple charge states after the electronic capture of  $^{57}\text{Co}$ .

$\text{CoCl}_2$  presents rhomboedral structure, of type  $\text{CdCl}_2$  (space group  $D_{3d}-3m$ ). The  $\text{Co}^{+2}$  ions stay in a nearly octahedral environment of  $6\text{Cl}^-$  anions, distorted along its trigonal (1,1,1) axe, which corresponds to the crystal c axe. Cations and anions are arranged in hexagonal layers perpendicular to the c axe, with the  $\text{Co}^{+2}$  layers being separated by two layers of  $\text{Cl}^-$  anions.

#### EXPERIMENTAL

The Mössbauer spectrometer was of constant acceleration type. The moving  $^{57}\text{Fe}$  enriched ferrocyanide absorber was kept at room temperature. The  $^{57}\text{Co}-\text{CoCl}_2$  source was prepared by evaporation of a solution of  $\text{CoCl}_2$  containing 0.5 to 1 mci of  $^{57}\text{Co}$ , followed by careful drying under  $\text{HCl}$  or  $\text{N}_2$  atmosphere at  $180^\circ\text{C}$ . Special precautions were taken against rehydration of the salt during the experiments. The Mössbauer parameters were obtained by a least square computer curve fitting, supposing that individual lines present theoretical Lorentzian line shape.

EXPERIMENTAL RESULTS

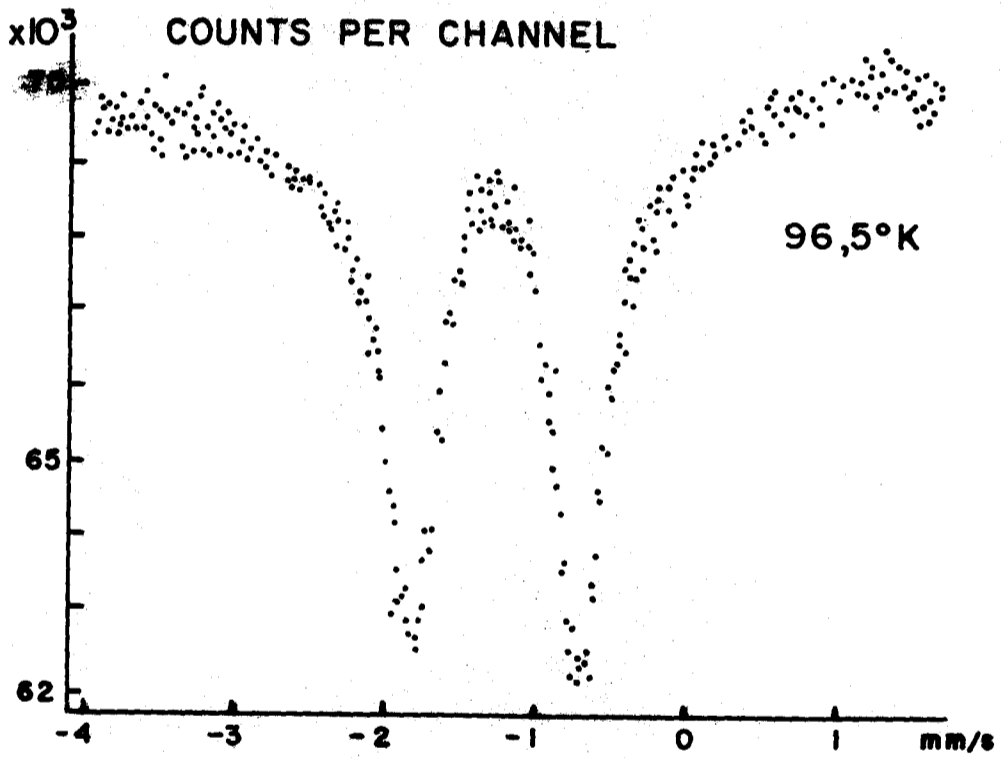
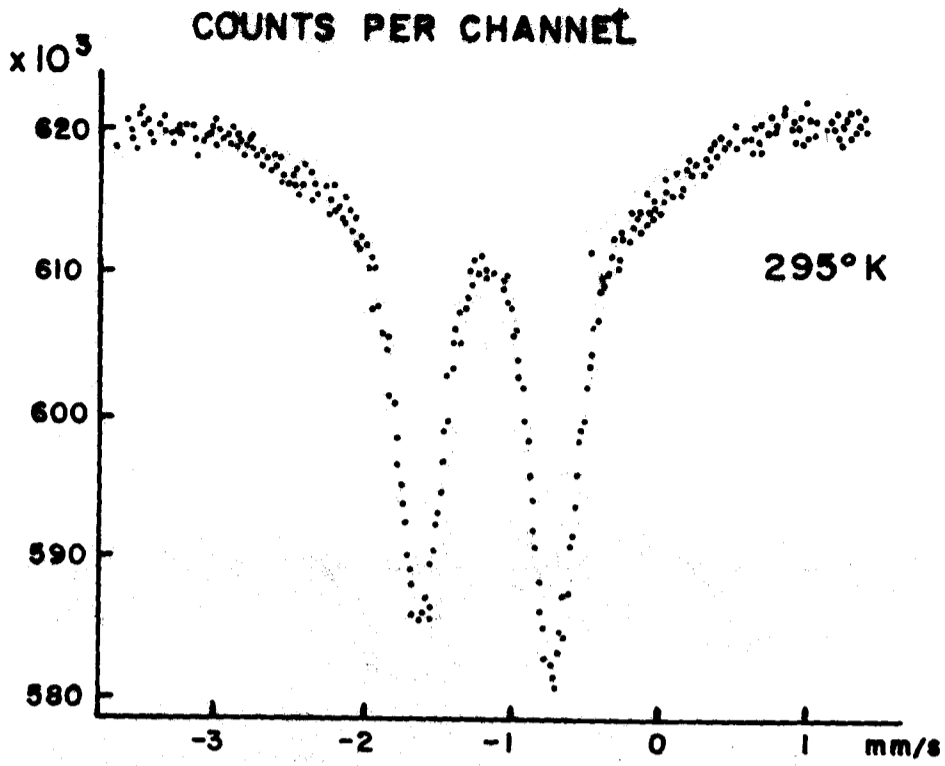
The Mössbauer spectra of  $^{57}\text{Co}-\text{CoCl}_2$  (figure) present two peaks of unequal heights and same linewidths. The isomer shift characterises ferrous ions; the quadrupole splitting is unusually low (typical values of  $\varepsilon$  for  $\text{Fe}^{+2}$  range from 2 to 3.5 mm/s) and strongly temperature dependant (Table). These features are however well in agreement with Mössbauer measurements of absorbers of ferrous halides <sup>2</sup>.

T	$-\sigma$	$\varepsilon$	$\Gamma$	$h_1/h_2$
95	1.27	1.02	0.41	1.06
295	1.14	0.92	0.38	1.11
375	1.05	0.81	0.37	1.19
470	1.00	0.69	0.45	1.23

Table: Mössbauer parameters of  $^{57}\text{Co}$  doped  $\text{CoCl}_2$  source. Temperature in °K, isomer shift:  $\sigma$  against  $\text{Na}_4\text{Fe}(\text{CN})_6 \cdot 10\text{H}_2\text{O}$ , quadrupole splitting:  $\varepsilon$  and linewidth:  $\Gamma$  all in mm/s.  $h_1/h_2$  is the ratio of the intensities of the two lines. The error upon  $\sigma$ ,  $\varepsilon$ , and  $\Gamma$  is 0,03 mm/s.

DISCUSSION OF THE EXPERIMENTAL RESULTS

The isomer shift of the ferrous ions formed after electronic capture of cobalt 57 in  $\text{CoCl}_2$  is higher than in ferrous chloride. This is attributed to a stronger delocalization of the 3d electrons, which might be induced by a pressure effect. (The ionic radius of  $\text{Fe}^{+2}$  is larger than that of  $\text{Co}^{+2}$ ) or be caused by an incomplete lattice rearrangement after the nuclear



Mössbauer spectra of  $^{57}\text{Co-CoCl}_2$

event. It has already been noticed that the isomer shifts of the ions thus formed by  $^{57}\text{Co}$  E. C. always imply a higher ionicity of the bonding than in the corresponding iron compounds <sup>3</sup>. The temperature variation is due to the second order Doppler shift ( $6 \times 10^{-4}$  mm/s/°K).

The quadrupole splitting is well described using earlier theoretical works by Ingalls and Gibb <sup>4</sup>. In the case of  $^{57}\text{Fe}-\text{CoCl}_2$ , the crystal field perturbation hamiltonian acting upon the  $^5\text{D}$  state is written:

$$V = V_{\text{cubic}} + V_{\text{trigonal}} = C_4(x^4 + y^4 + z^4 - \frac{3}{5}r^4) + B_2^0(3z^2 - r^2)$$

The cubic term splits the  $^5\text{D}$  state into an orbital doublet  $d_\gamma$  and a triplet  $d_\epsilon$ , whose energy separation is about  $7000 \text{ cm}^{-1}$  at  $295^\circ\text{K}$  <sup>5</sup>. The three-fold degeneracy of the  $d_\epsilon$  lower state is lifted by the trigonal field and the spin orbit interaction, whose effects are described by the equivalent operator hamiltonian  $-\lambda \hat{L} \hat{S} - \frac{\Delta_t}{3}(3 \hat{L}_z^2 - \hat{L}^2)$  acting upon the wave functions

$$\begin{aligned} & \frac{2}{3} \psi_{-2} + \frac{1}{3} \psi_1 \\ & \psi_0 \\ & \frac{1}{3} \psi_{-1} - \frac{2}{3} \psi_{-2} \end{aligned}$$

The  $15 \times 15$  perturbation matrix has been solved by several authors <sup>4</sup>. It may be concluded from the small value of  $\epsilon$  and from a comparison with isomorphous  $\text{FeCl}_2$  experiments that  $\Delta_t$  is positive.

Only  $q_{zz}$  contributes to the quadrupolar interaction since the crystal field presents axial symmetry. As the thermal transition frequencies between the sublevels are much higher than the quadrupole precession frequency ( $\approx 10^9 \text{ s}^{-1}$ ), the principal component of the electric field gradient is an average over the electronic states:

$$q_{zz} = \frac{\sum_n (q_{zz})_n e^{-\epsilon_n/kT}}{\sum_n e^{-\epsilon_n/kT}}$$

where  $(q_{zz})_n = \frac{4}{7} e \langle r^{-3} \rangle \langle \phi_n | 3\hat{l}_z^2 - \hat{l}^2 | \phi_n \rangle$  and  $\phi_n$  is the eigenvector of the  $n^{\text{th}}$  eigenstate with energy  $\epsilon_n$ . Neglecting the lattice contribution, the quadrupole splitting is then written in terms of a dimensionless factor  $F$  which is a function of  $\Delta_t$  and  $\lambda$ :

$$\epsilon = \frac{e q_{zz} Q}{2} = \frac{2}{7} e^2 Q(1-R) \alpha^2 \langle r_0^{-3} \rangle f(\Delta_t, \alpha^2 \lambda_0, T)$$

where  $Q$  is the quadrupole moment of the first excited nuclear state,  $R$  the Sternheimer polarization factor (+ 0.32),  $\alpha^2 = \frac{\lambda}{\lambda_0}$  the covalency factor accounting for the reduction of the spin orbit coupling constant  $\lambda$  and  $\langle r_0^{-3} \rangle$  from the free ion values  $\lambda_0$  (103  $\text{cm}^{-1}$ ) and  $\langle r_0^{-3} \rangle$  (4.8 a.u.).

Using Ingalls estimate for  $Q$  (0.29 barn) it is found that  $\epsilon = 5.62 \alpha^2 F(\Delta_t, \alpha^2 \lambda_0, T)$ . We supposed that the  $\lambda$  value is the same as in  $\text{FeCl}_2$ : 95  $\text{cm}^{-1}$  (2) and found  $\Delta_t = 160 \text{ cm}^{-1}$  with the help of the graphical results from reference 4.

However, the value of  $Q$  is not yet definitely established. Using  $Q = 0,18$  barns <sup>6</sup> and the same values for the other parameters, it is found that  $\Delta_t = 285 \text{ cm}^{-1}$ . The values of  $\Delta_t$  deduced from Mössbauer spectra are thus rather imprecise.

Furthermore, as is indicated by the isomer shift value, the spin orbit coupling constant should probably be somewhat higher for  $^{57}\text{Fe}^{+2}$  in  $\text{CoCl}_2$  than in  $\text{FeCl}_2$ , and the value of  $\lambda$  proposed in several works for  $\text{Fe}^{+2}$  in  $\text{FeCl}_2$  (2) seems to be anomalously high considering the covalency of the bonding. The main point is that the small value of the quadrupole interaction energy and its strong temperature dependence may well be explained by a strong mixing of the different electronic states under influence of the crystal field and spin orbit coupling and a small energy gap between the lowest levels.

The slight inequality of the peak heights of the spectra of  $\text{Fe}^{+2}$  formed in  $\text{CoCl}_2$  is attributed to an anisotropy of the Lamb Mössbauer factor <sup>7</sup>, in accordance with the layer structure of the compound. Relaxation effects and preferred orientation of the microcrystals were ruled out respectively by considering that both peaks have the same linewidth and that their intensity ratio does not vary with the orientation of the source relative to the  $\gamma$  ray axis. The increase of the spectra asymmetry with temperature is expected for the Goldanskii-Karyagin effect, which is invoked here.



CHARGE STATES OF  $^{57}\text{Fe}$  AFTER ELECTRON CAPTURE IN  $^{57}\text{Co}$ 

The interesting fact is that  $^{57}\text{Fe}$  is only found in the ferrous state in anhydrous  $\text{CoCl}_2$ . The difference against the hydrates of cobaltous chloride has already been discussed, and the appearance of ferric ions in the hydrates has been attributed to radiochemical reactions following the local autoradiolysis <sup>8</sup>.

Whereas only  $\text{Fe}^{+2}$  is observed in  $\text{CoCl}_2$ , the presence of  $\text{Fe}^{+2}$  and  $\text{Fe}^{+3}$  has been detected in  $\text{CoF}_2$  <sup>6</sup>. The appearance of ferric ions in the  $\text{CoF}_2$  lattice may be related to the increase of the lattice energy and of the Madelung constant on going from  $\text{CoCl}_2$  to  $\text{CoF}_2$ . Similar correlation has already been found in the complex fluorides  $\text{KCoF}_3$  and  $\text{K}_3\text{CoF}_6$  (3). This qualitative relation could mean that the higher charge states of  $^{57}\text{Fe}$  may be observed if the local variation of electrostatic energy associated with the change of charge is high enough against the ionization potential of the normal charge state of iron. This variation of electrostatic energy is of course related to the Madelung constant of the lattice. Calculations are under work to verify these hypotheses in several type of compounds.

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