# Cosmic crystallography: three multi-purpose functions 

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#### Abstract

A solid sphere is considered, with a uniformly distributed infinity of points. Two points being pseudorandomly chosen, the analytical probability density that their separation have a given value is computed, for three types of the underlying geometry: $E^{3}, H^{3}$ and $S^{3}$. Figures, graphs and histograms to complement this short note are given.


Key-Words: Cosmic crystallography; Curved three-spaces

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## 1 Introduction

The shape of the universe is presently under investigation, and cosmic crystallography (CC) is one of the various methods proposed to determine it [1].

The idea which supports the CC method is that if the universe is multiply connected then multiple images of a same cosmic object (a given quasar, say) may be observed in the sky. The separations between pairs of these images are correlated by the geometry and the topology of the spacetime; so if one selects a catalogue of observed images of cosmic objects and performs a histogram of the separations $l$ between the images, then these correlations manifest either as spikes (associated with Clifford translations) or as slight deformations of the histogram of the corresponding simply connected manifold [2].

A variant method was recently proposed by Fagundes and Gausmann [3], of subtracting from a histogram $\phi\left(l_{i}\right)$ of a multiply connected space a histogram of the corresponding simply connected space. The reported result was a plot with much oscillations in small scales.

In the present paper we propose an alternative to Fagundes-Gausmann method: we derive continuous probability density functions $\mathcal{F}(l)$ to be subtracted from the histogram $\phi\left(l_{i}\right)$, and thereby obtain a histogram sensibly more suitable for analysis.

In section 2 we derive the functions $\mathcal{F}(a, l)$ for the Euclidean, hyperbolic and elliptic geometries, and in section 3 we make a few comments.

## 2 Probability densities

We consider one of the simply connected spaces $E^{3}, H^{3}$ or $S^{3}$. In that space, a spherical solid ball $\mathcal{B}_{a}$ is taken, with radius $a$. The ball is assumed to contain an infinite number of pointlike objects, spatially distributed as uniformly as possible. We next select pseudorandomly two points of $\mathcal{B}_{a}$ and ask for the probability $\mathcal{F}(a, l) d l$ that the separation between the points lie between $l$ and $l+d l$.

## a. Euclidean geometry

In the Euclidean three-space take a ball $\mathcal{B}_{a}$ centred at the origin $O$ and select two points $P, Q$ in the ball; let $r \in[0, a]$ be the radial position of $P$ and let $l \leq 2 a$ be the distance from $P$ to $Q$ (see Figure 1).

Clearly the probability density $\mathcal{F}_{E}(a, r, l)$ of this configuration is proportional both to the area $\mathcal{S}_{E}(r)=4 \pi r^{2}$ of the locus of $P$ well as to the area of the locus of $Q$. When $r+l<a$ this locus is a sphere $S^{2}$ with area $\mathcal{S}_{E}(l)$, while when $r+l>a$ the locus is a spherical disk $D^{2}$ in $E^{3}$ with area $D_{E}(a, r, l)=(\pi l / r)\left[a^{2}-(l-r)^{2}\right]$.

The probability density of the configuration is then

$$
\begin{equation*}
\mathcal{F}_{E}(a, r, l)=k \mathcal{S}_{E}(r)\left\{\mathcal{S}_{E}(l) \times \Theta(a-l-r)+D_{E}(a, r, l) \times \Theta(l+r-a)\right\}, \tag{2.1}
\end{equation*}
$$

where $k$ is a normalization constant and $\Theta$ is the Heaviside function.
We integrate eq. (2.1) for $r \in[0, a]$, and finally obtain, for $l \in(0,2 a]$,

$$
\begin{equation*}
\mathcal{F}_{E}(a, l)=\frac{3 l^{2}}{16 a^{6}}(2 a-l)^{2}(l+4 a), \tag{2.2}
\end{equation*}
$$

where the value $k=9 /\left(16 \pi^{2} a^{6}\right)$ was set to satisfy the normalization condition

$$
\begin{equation*}
\int_{0}^{2 a} \mathcal{F}_{E}(a, l) d l=1 \tag{2.3}
\end{equation*}
$$

One often encounters in the literature reference to the probability density $\mathcal{P}(s)$ that the squared separation be $s$; since $s=l^{2}$ and $\mathcal{P}(s) d s=\mathcal{F}(l) d l$, then

$$
\begin{equation*}
\mathcal{P}_{E}(a, s)=\frac{3 \sqrt{s}}{32 a^{6}}(2 a-\sqrt{s})^{2}(\sqrt{s}+4 a) \tag{2.4}
\end{equation*}
$$

valid for $s \in\left(0,4 a^{2}\right]$.
In Figure 2 we reproduce a typical mean pair separation histogram (MPSH) for pseudorandomly distributed objects in an Euclidean solid sphere $\mathcal{B}_{a}$ with arbitrary radius, together with the corresponding probability density $\mathcal{F}_{E}(a, l)$. It should be stressed that, differently from the hyperbolic and elliptical cases, the shape of the function $\mathcal{F}_{E}(a, l)$ does not depend on the value of the radius $a$.

## b. Hyperbolic geometry

To obtain the probability density $\mathcal{F}_{H}(a, l)$ for the hyperbolic geometry we follow the same lines as before. The area of a sphere with radius $r$ is now $S_{H}(r)=4 \pi R^{2} \sinh ^{2} r / R$, where $R$ is the radius of curvature of the geometry; without loss of generality we henceforth set $R=1$. On the other hand, the area of a spherical disk $D^{2}$ in $H^{3}$ is (see Figure 1)

$$
\begin{equation*}
D_{H}(a, r, l)=2 \pi \sinh l[\sinh l-\cosh l \operatorname{coth} r+\cosh a \operatorname{csch} r], \tag{2.5}
\end{equation*}
$$

to be considered whenever $r+l>a$.
The probability density of the configuration is then

$$
\begin{equation*}
\mathcal{F}_{H}(a, r, l)=k S_{H}(r)\left[S_{H}(l) \times \Theta(a-l-r)+D_{H}(a, r, l) \times \Theta(l+r-a)\right] \tag{2.6}
\end{equation*}
$$

which upon integration for $r \in[0, a]$ and normalization gives

$$
\begin{equation*}
\mathcal{F}_{H}(a, l)=\frac{8 \sinh ^{2} l}{(\sinh 2 a-2 a)^{2}}[\cosh a \operatorname{sech}(l / 2) \sinh (a-l / 2)-(a-l / 2)] \tag{2.7}
\end{equation*}
$$

valid for $l \in(0,2 a]$.
In Figure 3 we reproduce graphs of $\mathcal{F}_{H}(a, l)$ for three values of the radius $a$. For $a \ll 1$ the function tends to the Euclidean one given in Figure 2, as expected. For increasing values of $a$ the function shifts towards the large values of $l$, and for $a \gg 1$ a strong concentration of $\mathcal{F}_{H}(a, l)$ is found near the extreme value $l=2 a$.

Figure 4 shows a typical MPSH in the hyperbolic three-space, together with the corresponding probability density $\mathcal{F}_{H}(a, l)$.

## c. Elliptic geometry

The basic strategy to obtain the probability density $\mathcal{F}_{S}(a, l)$ is the same as before, and the calculations are similar whenever the diameter $2 a$ of the ball is less than the separation $\pi R$ between antipodal points in the three-sphere; we then find for $\mathcal{F}_{S}(a, l)$ the expression, valid for $l \in(0,2 a]$,

$$
\begin{equation*}
\mathcal{F}_{S}(a<\pi / 2, l)=\frac{8 \sin ^{2} l}{(2 a-\sin 2 a)^{2}}[(a-l / 2)-\cos a \sec (l / 2) \sin (a-l / 2)], \tag{2.8}
\end{equation*}
$$

where we have taken $R=1$ without loss of generality.

However, the cases where $a>\pi / 2$ are considerably trickier to deal with, due to the treacherous connectivity of the spherical space $S^{3}$ and the new requirement that $l$ must not exceed $\pi$. A much larger quantity of trivial integrations now comes into scene, before the following expression is eventually obtained:

$$
\begin{align*}
\mathcal{F}_{S}(a, l)= & \frac{8 \sin ^{2} l}{[2 a-\sin 2 a]^{2}}\{[2 a-\sin 2 a-\pi]+  \tag{2.9}\\
& \Theta(2 \pi-2 a-l) \times[\sin 2 a+\pi-a-l / 2-\cos a \sec (l / 2) \sin (a-l / 2)]\},
\end{align*}
$$

valid for all $a \in(0, \pi]$ and $l \in(0, \min (2 a, \pi)]$.
In Figure 5 four graphs of $\mathcal{F}_{S}(a, l)$ are shown, for different values of the radius $a$ of the ball. For $a$ increasing from 0 to $\pi$ the function continuously shifts towards the smaller values of $l$. In particular, when $a=\pi / 2$ we have

$$
\begin{equation*}
\mathcal{F}_{S}(\pi / 2, l)=\frac{4}{\pi}(1-l / \pi) \sin ^{2} l, \tag{2.10}
\end{equation*}
$$

while when $a=\pi$ we have the harmonic, symmetric probability density

$$
\begin{equation*}
\mathcal{F}_{S}(\pi, l)=\frac{2}{\pi} \sin ^{2} l . \tag{2.11}
\end{equation*}
$$

In Figure 6 a sample MPSH in the spherical space is reproduced, together with the corresponding probability density $\mathcal{F}_{S}(a, l)$.

## 3 Comments

It is perhaps worth clarifying the meaning of the functions $\mathcal{F}(a, r, l)$ : if we pseudorandomly choose two points $P$ and $Q$ in a solid sphere with radius $a$ (see Figure 1) then $\mathcal{F}(a, r, l) d r d l$ is the probability that $P$ lies between the radial positions $r$ and $r+d r$, times the probability that the separation from $Q$ to $P$ lies between the values $l$ and $l+d l$. The form of $\mathcal{F}(a, r, l)$ clearly depends on the geometry one is concerned with.

Each histogram in Figures 2, 4, and 6 has $m=100$ subintervals and is a MPSH - a mean pair separation histogram comprising $K=10$ comparable catalogues with $N=50$ objects each [2]. All computer-generated histograms assume a homogeneous distribution, as described by Lehoucq, Luminet and Uzan [4].

A close inspection of Figure 2 shows that the most probable separation between two arbitrarily chosen points in an Euclidean solid ball is slightly greater than the radius of the ball; also the maximum of $\mathcal{P}_{E}(a, s)$ in eq.(2.4) occurs when $s / 4 a^{2}=0.134$, in agreement with plots of Fagundes and Gausmann [5].

A characteristic feature of the hyperbolic geometries is that at large distances there is more space than in the Euclidean geometries; this fact is clearly exhibited in Figure 3, which shows a strong predominance of large separations $l$ when the radius $a$ of the solid sphere is large. It is worth noting that Fagundes and Gausmann [3] obtained histograms with $a=2.34$ which closely resemble ours with $a=2$ in Figure 4. In contrast, the hyperbolic histograms given by Lehoucq et al. [4] use radius $a$ nearly 0.6 , so they are similar to our Euclidean one.

Oppositely to the hyperbolic case, in distant places endowed with the elliptic geometry there is less space than in the Euclidean geometry; this is nicely illustrated in Figure 5. Indeed, with increasing $a$ ( increasing solid ball) the probability of finding small distances $l$ (in comparison with $a$ ) also increases. We further note that when $a$ increases from $\pi / 2$ to $\pi$ the ratio $l_{\text {max }} / 2 a$ recedes from 1 to $1 / 2$.

## References

[1] R.Lehoucq, M.Lachièze-Rey and J.-P.Luminet, Astron. Astroph. 313, 339-346 (1996).
[2] G.I.Gomero, A.F.F.Teixeira, M.J.Rebouças and A.Bernui, gr-qc/9811038 (1998).
[3] H.V.Fagundes and E.Gausmann, astro-ph/9811368 (1998).
[4] R.Lehoucq, J.-P.Luminet and J.-Ph.Uzan, astro-ph/9811107 (1998).
[5] H.V.Fagundes and E.Gausmann, Phys. Letters A 238, 235-238 (1998).

Figure 2: Figure 2. In the Euclidean space $E^{3}$, a sample MPSH for a solid ball with
arbitrary radius $a$, together with the corresponding probability density eq.(2.2).
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