The Non-Abelian Aharonov-Casher Effect in the Framework of Feynman Pseudo-Classical Path Integrals

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Abstract

I analyze the non-Abelian Aharonov-Casher effect in the Feynman path integral framework.

The Aharonov-Casher effect for electromagnetism has been discussed in the literature as a basic building to theory's of high T_e -superconductivity ([1],[2]). My aim in this note is to generalize this effect for the non-abelian case by making use of Feynman pseudoclassical path integrals as proposed by myself on refs. [3].

Let me start our study by considering the Grassmanian path integral representation of ref. [3] for the Green function of a U(N) neutral spin one half relativistic particle with a non zero coloured magnetic dipole moment $\bar{\mu}_m$ interacting with a constant U(N) non abelian field strenght $F_{\alpha\beta}$:

$$G_{(\mu,\eta)}^{(a,b)}[X_{\mu}^{(1)}, X_{\mu}^{(2)}] = \int_{0}^{\infty} ds \int_{\substack{X^{\mu}(0) = X_{\mu}^{(1)} \\ X^{\mu}(s) = X_{\mu}^{(2)}}} \prod [dx^{\mu}(\sigma)]$$

$$IP_{colour} IP_{spin} \left\{ exp \ i/4\bar{\mu} \int_{0}^{s} [\gamma^{\alpha}, \gamma^{\beta}](\sigma) F_{\alpha\beta}(X^{\gamma}(\sigma)) d\sigma \right\}_{(\mu,\eta)}^{(a,b)}$$
(1)

where I have evaluated exactly the Grassmanians path integrals associated to the colour charge and Lorentz spin degrees of freedom yielding, thus, the U(N)-colour and Lorentz group valued Mandelstan phase factor written in eq. (1), namelly

$$W^{F} = \prod_{\substack{0 < \sigma < s}} [d\theta(\sigma)][d\theta^{*}(\sigma)][d\psi_{\mu}(\sigma)](\theta_{a}(0)\theta^{*}_{b}(s))(\psi_{\mu}(0)\psi_{\eta}(s))$$

$$exp\left\{\frac{1}{2}i\int_{0}^{s}(\theta\theta^{*})(s)ds + \frac{1}{2}i\int_{0}^{s}\psi_{\mu}(s)\dot{\psi}_{\mu}(s)ds\right\}$$

$$exp\left\{\int_{0}^{s}d\sigma\frac{1}{2}[\psi_{\mu}(\sigma),\psi_{\eta}(\sigma)F^{i}_{\mu\eta}[X^{\alpha}(\sigma)(\theta_{a}(\lambda_{i})^{ab}\theta^{*}_{b})(\sigma)]\right\}$$
(2)

Let me analyse the following non abelian Aharonov-Chaser experiment ([3]): A quantum U(N) neutral particle with a U(N)-coloured magnetic moment travels along a closed trajectory on the region M where $F^a_{\alpha\beta}(X^{\rho})\lambda_a$ is a constant field strenght. We remark now that the only Feynman bosonic trajectory that contributes to the fermionic Mandelstan phase factor in eq. (1) comes from the classical trajectory $X^{CL}_{\mu}(\sigma)$. This result may obtained by expanding in power of \hbar , the fermionic Mandelstan phase factor eq. (2):

$$I\!\!P_{colour} I\!\!P_{spin} \left\{ exp \ i/\hbar\bar{\mu}_m \int_0^s [\gamma^{\alpha}, \gamma^{\beta}](\sigma) F_{\alpha\beta}(X_{\mu}^{c\ell}(\sigma) + \hbar X_{\mu}^q(\sigma)) d\sigma \right\} - \\
 I\!\!P_{colour} I\!\!P_{spin} \left\{ exp \ i/\hbar\bar{\mu}_m \int_0^s d\sigma [\gamma^{\alpha}, \gamma^{\beta}](\sigma) F_{\alpha\beta}(X_{\mu}^{c\ell}(\sigma)) \right\} = \\
 = \sum_{n=1}^{\infty} \frac{(\hbar)^n}{n!} \int_0^s dX_{\alpha_1}^q(\sigma_1) \cdots \int_0^s dX_{\alpha_n}^q(\sigma_n) \\
 \left[\frac{\delta^{(n)}}{\delta\eta_{\alpha_1}(\sigma_1) \cdots \delta\eta_{\alpha_n}(\sigma_n)} I\!\!P \left\{ exp \ i/\hbar\bar{\mu}_m \int_0^s [\gamma^{\alpha}, \gamma^{\beta}](\sigma) F_{\alpha\beta}(\eta_{\mu}(\sigma)) d\sigma \right\} \right]_{\eta_{\mu}(\sigma) = X_{\mu}^{CL}(\sigma)} (3)$$

We point out that functional derivatives in eq. (3) are multiplicatively proportional to derivatives of the field strenght, which vanish as $F_{\alpha\beta}(x)$ is taken to be a constant field strenght on M. We have, thus, the validity of the above mentioned remark. We get therefore that the wave function of a U(N) spin one half neutral particle with U(N) non zero magnetic dipole moment in the non abelian Aharonov-Casher experiment develops a topological phase given explicitly by the Fermionic Wilson loop of ref. ([3]) for a neutral particle

$$W^{F} = Tr_{spin}Tr_{colour} \left[I\!\!P_{colour} I\!\!P_{spin} \left\{ exp \ i/\hbar \ \bar{\mu}_{m} \int_{0}^{s} [\gamma^{\alpha}, \gamma^{\beta}](\sigma) F_{\alpha\beta}(X^{CL}_{\mu}(\sigma)) d\sigma \right\} \right]$$
(4)

This result confirms the explicit appearance of Lorentz spin degrees of freedom on the phase shift of the Aharonov-Casher wave function predicted in refs. [1], [2].

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